

## Chapter 18

# A First Look at Stochastic Integrals with the Wiener Process

Section 18.1 touches briefly on the martingale characterization of the Wiener process.

Section 18.2 gives a heuristic introduction to stochastic integrals, via Euler's method for approximating ordinary integrals.

### 18.1 Martingale Characterization of the Wiener Process

Because the Wiener process is a Lévy process (Example 139), it is self-similar in the sense of Definition 147. That is, for any  $a > 0$ ,  $W(at) \stackrel{d}{=} a^{1/2}W(t)$ . In fact, if we define a new process  $W_a$  through  $W_a(t, \omega) = a^{-1/2}W(at, \omega)$ , then  $W_a$  is itself a Wiener process. Thus the whole process is self-similar. This is only one of several sorts of self-similarities in the Wiener process. Another is sometimes called *spatial homogeneity*:  $W_\tau$ , defined through  $W_\tau(t, \omega) = W(t + \tau, \omega) - W(\tau, \omega)$  is also a Wiener process. That is, if we “re-zero” to the state of the Wiener process  $W(\tau)$  at an arbitrary time  $\tau$ , the new process looks just like the old process.  $-W(t)$ , obviously, is also a Wiener process.

Related to these properties is the fact that  $W^2(t) - t$  is a martingale with respect to  $\{\mathcal{F}_t^W\}$ . (This is easily shown with a little algebra.) What is more surprising is that this is enough to characterize the Wiener process.

**Theorem 224 (Martingale Characterization of the Wiener Process)** *If  $M(t)$  is a continuous martingale, and  $M^2(t) - t$  is also a martingale, then  $M(t)$  is a Wiener process.*

There are some very clean proofs of this theorem<sup>1</sup> — but they require us to use stochastic calculus! Doob (1953, pp. 384ff) gives a proof which does not, however. The details of his proof are messy, but the basic idea is to get the central limit theorem to apply, using the martingale property of  $M^2(t) - t$  to get the variance to grow linearly with time and to get independent increments, and then seeing that between any two times  $t_1$  and  $t_2$ , we can fit arbitrarily many little increments so we can use the CLT.

We will return to this result as an illustration of the stochastic calculus (Theorem 247).

## 18.2 A Heuristic Introduction to Stochastic Integrals

Euler's method is perhaps the most basic method for numerically approximating integrals. If we want to evaluate  $I(x) \equiv \int_a^b x(t)dt$ , then we pick  $n$  intervals of time, with boundaries  $a = t_0 < t_1 < \dots < t_n = b$ , and set

$$I_n(x) = \sum_{i=1}^n x(t_{i-1})(t_i - t_{i-1})$$

Then  $I_n(x) \rightarrow I(x)$ , if  $x$  is well-behaved and the length of the largest interval  $\rightarrow 0$ . If we want to evaluate  $\int_{t=a}^{t=b} x(t)dw$ , where  $w$  is another function of  $t$ , the natural thing to do is to get the derivative of  $w$ ,  $w'$ , replace the integrand by  $x(t)w'(t)$ , and perform the integral with respect to  $t$ . The approximating sums are then

$$\sum_{i=1}^n x(t_{i-1})w'(t_{i-1})(t_i - t_{i-1}) \quad (18.1)$$

Alternately, we could, if  $w(t)$  is nice enough, approximate the integral by

$$\sum_{i=1}^n x(t_{i-1})(w(t_i) - w(t_{i-1})) \quad (18.2)$$

even if  $w'$  doesn't exist.

(You may be more familiar with using Euler's method to solve ODEs,  $dx/dt = f(x)$ . Then one generally picks a  $\Delta t$ , and iterates

$$x(t + \Delta t) = x(t) + f(x)\Delta t \quad (18.3)$$

from the initial condition  $x(t_0) = x_0$ , and uses linear interpolation to get a continuous, almost-everywhere-differentiable curve. Remarkably enough, this converges on the actual solution as  $\Delta t$  shrinks (Arnol'd, 1973).)

Let's try to carry all this over to random functions of time  $X(t)$  and  $W(t)$ . The integral  $\int X(t)dt$  is generally not a problem — we just find a version of  $X$

<sup>1</sup>See especially Ethier and Kurtz (1986, Theorem 5.2.12, p. 290).

with measurable sample paths (Section 8.2).  $\int X(t)dW$  is also comprehensible if  $dW/dt$  exists (almost surely). Unfortunately, we've seen that this is not the case for the Wiener process, which (as you can tell from the  $W$ ) is what we'd really like to use here. So we can't approximate the integral with a sum like Eq. 18.1. But there's nothing preventing us from using one like Eq. 18.2, since that only demands increments of  $W$ . So what we would like to say is that

$$\int_{t=a}^{t=b} X(t)dW \equiv \lim_{n \rightarrow \infty} \sum_{i=1}^n X(t_{i-1})(W(t_i) - W(t_{i-1})) \quad (18.4)$$

This is a crude-but-workable approach to numerically evaluating stochastic integrals, and apparently how the first stochastic integrals were defined, back in the 1920s. Notice that it is going to make the integral a *random variable*, i.e., a measurable function of  $\omega$ . Notice also that I haven't said anything yet which should lead you to believe that the limit on the right-hand side exists, in any sense, or that it is independent of the choice of partitions  $a = t_0 < t_1 < \dots < t_n = b$ . The next chapter will attempt to rectify this.

(When it comes to the SDE  $dX = f(X)dt + g(X)dW$ , the counterpart of Eq. 18.3 is

$$X(t + \Delta t) = X(t) + f(X(t))\Delta t + g(X(t))\Delta W \quad (18.5)$$

where  $\Delta W = W(t + \Delta t) - W(t)$ , and again we use linear interpolation in between the points, starting from  $X(t_0) = x_0$ .)